Decoherence Suppression of a Solid State Qubit Using Uncollapsing

Kyle Michael Keane

Alexander N. Korotkov

University of California, Riverside

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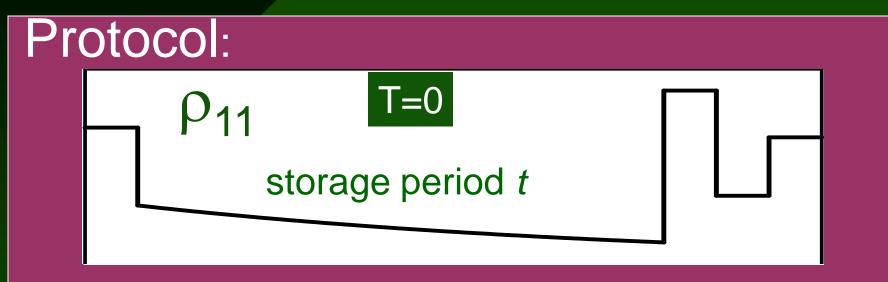


Other Proposed Methods

- Quantum Error Correcting Codes
 Complicated and Resource Hungry
- Decoherence-Free Subspace
 Complicated and Resource Hungry
- Dynamical Decoupling

 Does Not Protect Against Markovian Processes

Our Method: Phase Qubit



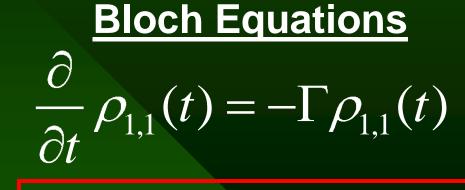


almost the same as existing experiment!

Phase Qubit: <u>Partial Measurement (PM)</u>

 $|1\rangle$

 $|0\rangle$



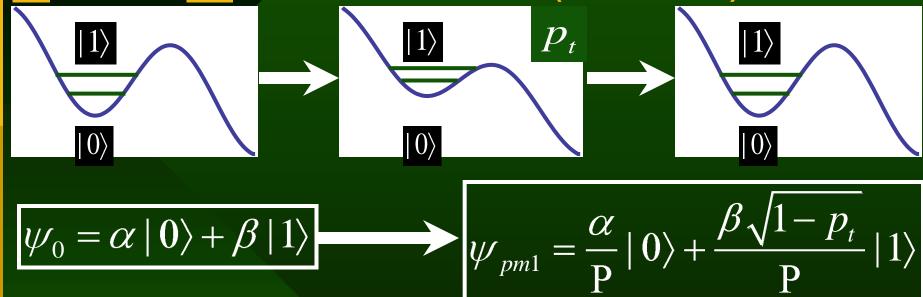
 $|1\rangle$

 $|0\rangle$

$$P_{T^{C}|1}(t) = e^{-\Gamma t} \equiv 1 - p_t(t)$$

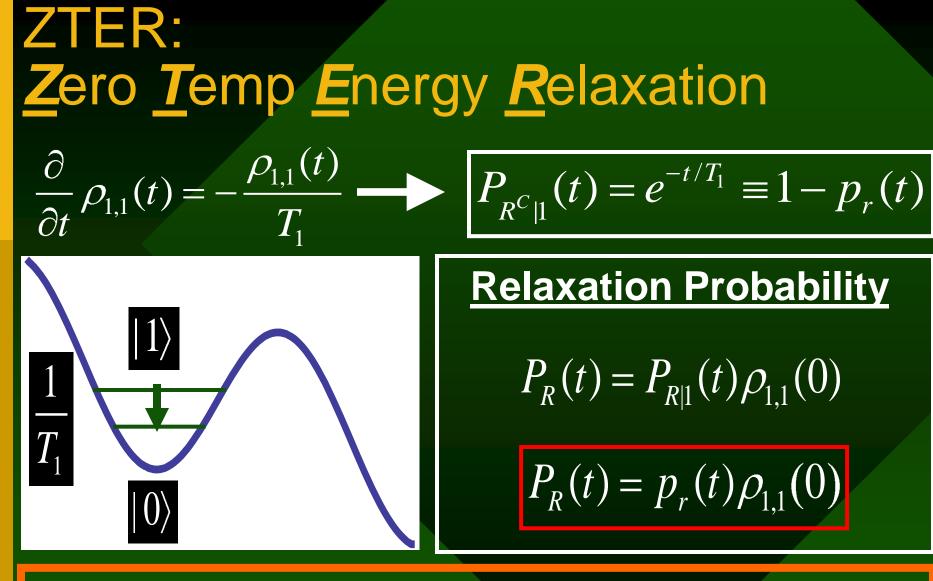
 $P_{T^{C}}(t) = \rho_{0,0}(0) + [1 - p_{t}(t)]\rho_{1,1}(0)$

Phase Qubit: Effect of <u>Null</u> <u>Result</u> <u>Partial</u> <u>Measurement</u> (NRPM)

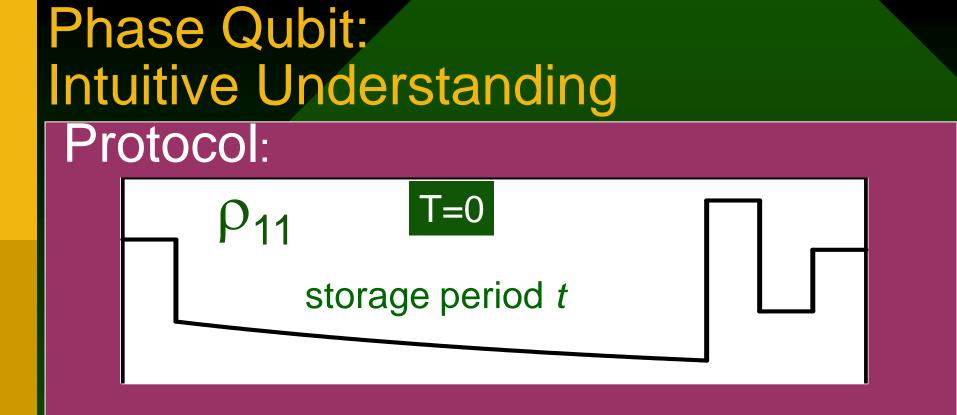


The NRPM reduces the probability that the qubit is in state $|1\rangle$ after the NRPM is complete

$$\frac{\alpha}{P} > \alpha \qquad \frac{\beta \sqrt{1-p_t}}{P} < \beta$$
$$P \equiv \sqrt{|\alpha|^2 + |\beta|^2 (1-p)}$$



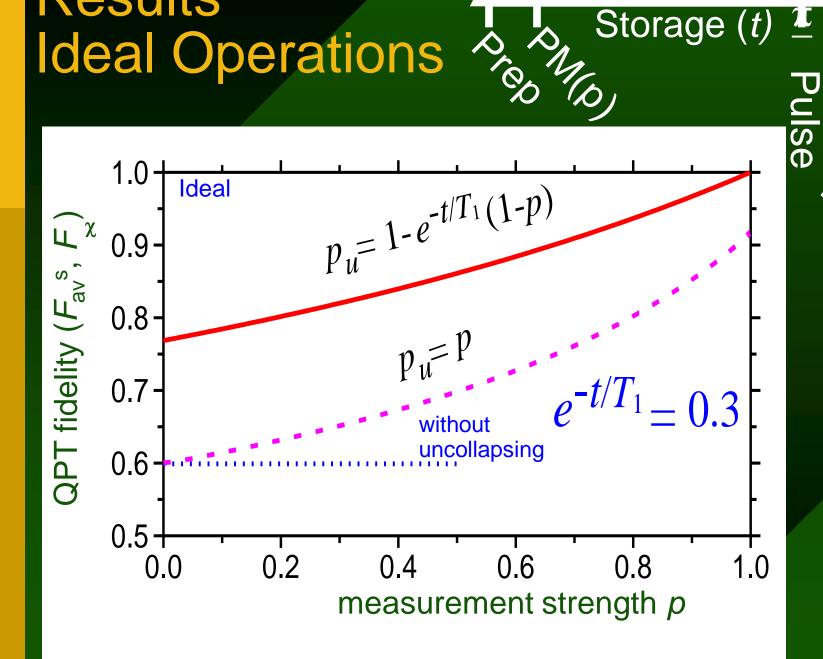
The probability that a qubit relaxes during a ZTER period is proportional to the probability that it was in state $|1\rangle$ at the beginning of the ZTER period



The first PM reduces the probability that the qubit is in state $|1\rangle$ after the PM is complete

The probability that the qubit relaxes during the storage period is proportional to the probability that it was in state $|1\rangle$ at the beginning of the storage

Results **Ideal Operations**



π

Pulse

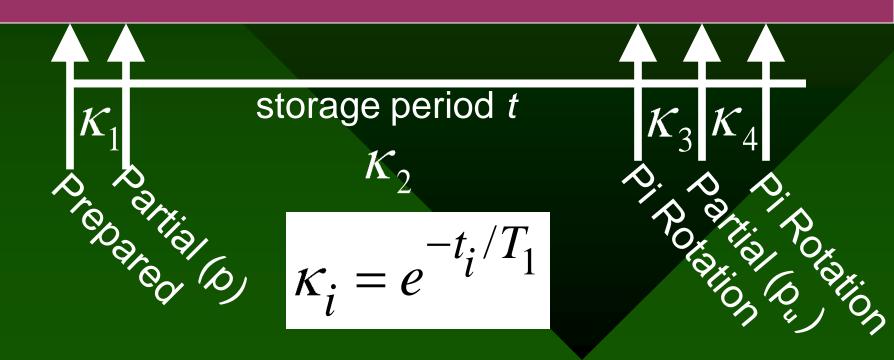
110

Pu

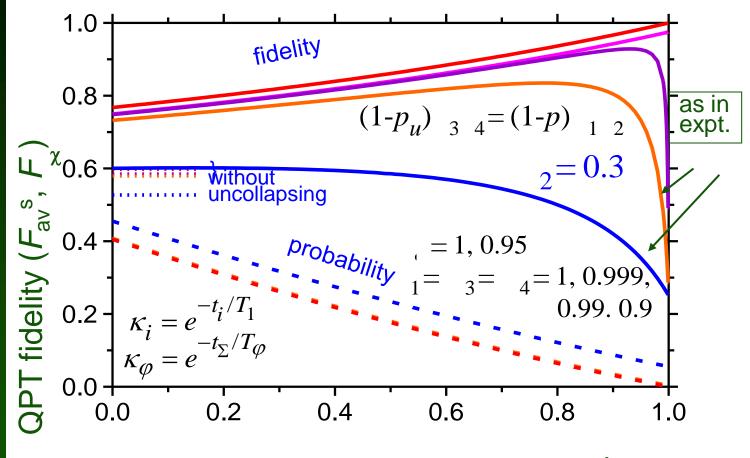
Storage (t)

Non-Ideal Operations Protocol:





Results Non-Ideal Operations



measurement strength p

Conclusions

Decoherence Suppression by Uncollapsing:

- Requires <u>NO</u> Extra Resources
- Protects Against Markovian Processes
- Works Even With Non-Ideal Operations
- Can Be Experimentally Realized With a Very Simple Extension of a Previous Experiment